

Contents lists available at ScienceDirect

Journal of Solid State Chemistry



journal homepage: www.elsevier.com/locate/jssc

Phase diagram, crystal chemistry and thermoelectric properties of compounds in the Ca–Co–Zn–O system

W. Wong-Ng^{a,*}, T. Luo^c, W. Xie^d, W.H. Tang^e, J.A. Kaduk^f, Q. Huang^b, Y. Yan^a, S. Chattopadhyay^f, X. Tang^g, T. Tritt^d

^a Ceramics Division, NIST, Gaithersburg, MD 20899, USA

^b NIST Center for Neutron Research, USA

^c Chemistry and Biochemistry Department, University of Maryland, MD 20745, USA

^d Physics Dept., Clemson University, Clemson, SC 29631, USA

^e Department of Physics, Zhejiang Sci-Tech University, Hangzhou 310018, China

f BCPS Dept., Illinois Institute of Technology, Chicago, IL 60616, USA

^g Physics Department, Wuhan University, Wuhan 430070, China

ARTICLE INFO

Article history: Received 29 March 2011 Received in revised form 10 June 2011 Accepted 13 June 2011 Available online 21 June 2011

Keywords: Ca-Co-Zn-O phase equilibria Structure of Ca₃(Co_{1.95}Zn_{0.05})O_{6-z} Structure of (Co_{2.7}Zn_{0.3})O_{4-z} Thermoelectric properties of Ca₃(CoZn)₄O_{9-z} and Ca₃(CoZn)₂O_{6-z}

ABSTRACT

In the Ca–Co–Zn–O system, we have determined the tie-line relationships and the thermoelectric properties, solid solution limits, and structures of two low-dimensional cobaltite series, Ca₃(Co, Zn)₄O_{9-z} and Ca₃(Co,Zn)₂O_{6-z} at 885 °C in air. In Ca₃(Co,Zn)₄O_{9-z}, which has a misfit layered structure, Zn was found to substitute in the Co site to a limit of Ca₃(Co,3,8Zn_{0.2})O_{9-z}. The compound Ca₃(Co,Zn)₂O_{6-z} (*n*=1 member of the homologous series, Ca_{*n*+2}(Co,Z_{*n*})_{*n*}(Co,Zn)'O_{3*n*+3-z}) consists of one-dimensional parallel (Co,Zn)₂O₆⁶⁻ chains that are built from successive alternating face-sharing (Co,Zn)O₆ trigonal prisms and 'n' units of (Co,Zn)O₆ octahedra along the hexagonal *c*-axis. Zn substitutes in the Co site of Ca₃Co₂O₆ to a small amount of approximately Ca₃(Co_{1.95}Zn_{0.05})O_{6-z}. In the ZnO-CoO_z system, Zn substitutes in the tetrahedral Co site of Ca₃O₄ to the maximum amount of (Co_{2.4}Zn_{0.3})O_{4-z}, (Zn_{0.94}Co_{0.06})O, and Ca₃(Co_{1.95}Zn_{0.05})O_{6-z} are described. Despite the Ca₃(Co, Zn)₂O_{6-z} series having reasonably high Seebeck coefficients and relatively low thermal conductivity, the electrical resistivity values of its members are too high to achieve high figure of merit, *ZT*.

© 2011 Elsevier Inc. All rights reserved.

1. Introduction

Over the last decade, the increase in global interest in research and development on thermoelectric (TE) materials has been partly due to the soaring short-term demand for energy and partly due to the need to create a sustainable energy future. The efficiency and performance of thermoelectric energy conversion is related to the dimensionless figure of merit (*ZT*) of the thermoelectric materials, given by $ZT=S^2\sigma T/\kappa$, where *T* is the absolute temperature, *S* is the Seebeck coefficient or thermoelectric power, σ is the electrical conductivity, and *k* is the thermal conductivity [1]. Only a small number of materials have been found to have practical industrial applications and they all have *ZT* values around or below 1.0. Optimization of the *ZT* values is not a straight-forward process because *S*, σ , and κ are

* Corresponding author *E-mail address:* Winnie.wong-ng@nist.gov (W. Wong-Ng). interrelated. Recent reports that relatively high *ZT* values are possible in both thin film and bulk forms have revitalized interest in thermoelectric materials development [2–7].

The stability of metal oxides at high temperature has made these materials highly relevant to waste heat conversion applications. For example, the low-dimensional cobaltites that include $NaCoO_z$ [8], $Ca_2Co_3O_6$ [9,10], and $Ca_3Co_4O_9$ [11–14] exhibit the coexistence of large Seebeck coefficient and relatively low thermal conductivity. Consequently in recent years, considerable research has been conducted on oxides for thermoelectric applications.

Phase equilibrium diagrams provide road maps for processing and facilitating an understanding of material properties. Because of the promising properties of calcium-containing cobaltite, zincdoped calcium cobaltites may also offer desirable properties. This paper discusses the phase compatibility relationships of the Ca–Co–O system, and crystal chemistry, crystallography, and thermoelectric property measurements of selected series of compounds in this system, including Ca₃(Co, Zn)₄O_{9–z}, and the n=1 Zn-doped homologous series, Ca_{n+2}(Co,Zn)_n(Co,Zn)'O_{3<math>n+3-z}.</sub>

^{0022-4596/\$ -} see front matter \circledcirc 2011 Elsevier Inc. All rights reserved. doi:10.1016/j.jssc.2011.06.013

2. Experimental¹

2.1. Sample preparation

Using high temperature solid-state synthesis techniques, 47 samples were prepared from stoichiometric amounts of CaCO₃, Co₃O₄, and ZnO (Table 1). These samples were mixed, pelletized, and annealed at 750 °C for one day, and subsequently annealed at 885 °C with intermediate grindings and pelletizations for another 7 days.

Powder X-ray diffraction was used to investigate the phase purity and establish phase relationships. These experiments were carried out using a Phillips X-ray powder diffractometer with CuK α radiation and equipped with a series of Soller slits and a scintillation counter. The 2θ scanning range was from 10° to 80°, and the step interval was 0.03°. The ICDD PDF reference diffraction patterns of the Ca–Co–Zn–O systems [15] were used for performing phase identification.

For property measurements, the spark plasma sintering (SPS) technique was used for the preparation of high density pellets. The raw powder materials were consolidated into densified pellets at 1000 °C for 5 min under a pressure of 30 MPa. All samples have 90% or higher relative densities. The pellets were cut into bars of approximate dimensions of $2\text{mm} \times 2 \text{ mm} \times 10 \text{ mm}$ for the measurement of properties.

2.2. Structural studies using synchroton and neutron diffraction

2.2.1. Synchroton scattering studies

To study the structure of selected phases in the Ca–Co–Zn–O system (in particular, Ca₃(Co_{1.95}Zn_{0.05})₂O_{6–z} and (Zn_{0.98}Co_{0.02})O), high resolution synchrotron powder diffraction data were collected using beamline 11-BM at the Advanced Photon Source (APS), Argonne National Laboratory, using an average wavelength of 0.412253 Å. Discrete detectors covering an angular range from 6° to 16° 2 θ are scanned over a 34° 2 θ range, with data points collected every 0.001° 2 θ and a scan speed of 0.01°/s.

The 11-BM instrument uses X-ray optics with two platinumstriped mirrors and a double-crystal Si(111) monochromator, where the second crystal has an adjustable sagittal bend [16]. Ion chambers monitor incident flux. A vertical Huber 480 goniometer, equipped with a Heidenhain encoder, positions an analyzer system comprised of twelve perfect Si(111) analyzers and twelve Oxford-Danfysik LaCl₃ scintillators, with a spacing of 2° [17]. Analyzer orientation can be adjusted individually on two axes. A three-axis translation stage holds the sample mounting and allows it to be spun, typically at \approx 5400 rpm (90 Hz). A Mitsubishi robotic arm is used to mount and dismount samples on the diffractometer [18].

The diffractometer is controlled via EPICS [19]. Data are collected while continually scanning the diffractometer 2θ arm. A mixture of NIST standard reference materials, Si (SRM 640c) and Al₂O₃ (SRM 676) [20] is used to calibrate the instrument, where the Si lattice constant determines the wavelength for each detector. Corrections are applied for detector sensitivity, 2θ offset, small differences in wavelength between detectors, and the source intensity, as noted by the ion chamber before merging the data into a single set of intensities evenly spaced in 2θ .

All data processing and structural refinements were carried out using the Rietveld refinement technique with the GSAS Suite [21,22]. In general, included in the refinements were the atomic

Та	Ы	e	1	
1 u		•		

Samples prepared for the phase equilibria study of the Ca–Co–Zn–O system at 885 $^{\circ}$ C in air. The values given are the relative mol% of metallic components.

	Ca	Со	Zn
1	42.86	53.57	3.57
2	42.86	50.00	7.14
3	42.86	46.43	10.71
4	60	36	4
5	60	32	8
6	32.5	32.5	35
7	60	20	20
8	0	66.67	33.33
9	70	20	10
10	10	75	15
11	42.86	57	0.14
12	42.86	56.43	0.71
13	42.86	55.71	1.43
14	60	39.8	0.2
15	60	39	1
16	60	38	2
17	50	0	50
18	99.5	0	0.5
19	97.5	0	2.5
20	0.5	0	99.5
21	2.5	0	97.5
22	0	95	5
23	0	90	10
24	0	5	95
25	0	10	90
26	0	83.33	16.67
27	0	80	20
28	0	76.67	23.33
29	60	37	3
30	0	3	97
31	0	4	96
32	75	24	1
33	50	48	2
34	20	75	5
35	20	1	79
36	0	2	98
37	60	40	0
38	42.86	57.14	0
39	37	3	60
40	28	18	54
41	38	25	37
42	30	37.5	32.5
43	30	40	30
44	15	55	30
45	0	6.67	93.33
46	0	8.33	91.67
47	37.5	25	37.5

coordinates, anisotropic metal displacement coefficients, isotropic oxygen displacement coefficients, a scale factor, and the lattice parameters. The peak profiles were described using a pseudo-Voigt function; the Gaussian *U* and *W*, the Cauchy *X* and anisotropic strain terms were refined. The backgrounds were described using a 6-term shifted Chebyshev polynomial.

X-ray powder reference patterns for samples in the Ca–Co–Zn–O system were obtained using a Rietveld pattern decomposition technique. These patterns represent ideal specimen patterns. They are corrected for systematic errors both in *d*-spacing and intensity. The reported peak positions are calculated from the refined lattice parameters, as this represents the best measure of the true positions.

2.2.2. Neutron diffraction studies

Neutron powder diffraction (NPD) data for $Zn_{0.3}Co_{2.7}O_{4-z}$ were collected at the NIST Center for Neutron Research on the high resolution powder neutron diffractometer (BT1), with monochromatic neutrons of wavelength 1.5403 Å produced by

¹ The purpose of identifying the equipment in this article is to specify the experimental procedure. Such identification does not imply recommendation or endorsement by the National Institute of Standards and Technology.

a Cu(311) monochromator. Collimations before and after the monochromator and after the sample were 15', 20', and 7' full-width-at-half-maximum (FWHM), respectively. Data were collected in the 2θ range of 3–168° with a step size of 0.05° at room temperature. Refinement of the nuclear structure was carried out using the neutron powder diffraction data and the Rietveld refinement technique using the software GSAS [21,22]. The neutron scattering amplitudes used in the refinements are 0.568, 0.253, and 0.581 for Zn, Co, and O, respectively.

2.2.3. EXAFS studies

EXAFS was used to evaluate the exact site substitution of Zn on the Co site in compounds where Zn substitution is only to a small extent. EXAFS measurements were performed on both Zn edge and Co edges of all the samples at the MRCAT 10-ID beamline located at the Advanced Photon Source in Argonne National Laboratory. Apart from the samples, we measured a solution of zinc nitrate in water (where zinc is 6-fold coordinated) [23]. ZnO powder was also measured since Zn in this compound is tetrahedrally bonded to oxygen atoms. With every EXAFS measurement, the zinc metal foil was measured as reference. The Zn edge of the sample of interest is placed along with Zn edges from Zn foil, ZnO powder, and $(Zn(NO_3)_2$ solution. The data were processed using Athena [24] and the data was modeled with Artemis [25] to find out the percentage of Zn atoms bonded tetrahedrally and octahedrally.

2.3. Thermoelectric property measurements

The thermal diffusivity was measured by the laser flash method using the Netzsch LFA457 system. Specific heat was determined by differential scanning calorimetry using a Netzsch DSC-404C. The density was calculated from the chemical formula and the unit cell volume. The resulting high temperature thermal conductivity, κ , was calculated from the measured thermal diffusivity, D, specific heat, C_P , and density, ρ_D , from the relationship: $\kappa = DC_P \rho_D$. The bulk materials were then cut into rectangular bars with the approximate dimensions of 2 mm \times 2 mm \times 10 mm. These bars were measured for electrical conductivity and Seebeck coefficient by a commercial ZEM-2 unit (Ulvac Riko, Inc.) under an inert gas (He) atmosphere from 300 to 900 K. Uncertainties in the electrical conductivity and thermal conductivity values are on the order of 5-7% due in large part in determination of the sample dimensions. The thermopower is known within about 2% and thermal diffusivity to within about 5%.

3. Results and discussion

Fig. 1 gives the phase diagram of the Ca–Co–Zn–O system that was determined at 885 °C in air. The phase relationships between solid solutions and other phases are expressed as tie-line bundles. The crystal chemistry, crystallography, and thermoelectric properties of various phases in the binary as well as the ternary oxide systems are discussed below.

3.1. CaO-ZnO system

No binary compound was found in the CaO–ZnO system. The literature reported that the CaZnO₂ phase was only successfully prepared under different processing conditions [15]. CaZnO₂ was also confirmed to be absent in a number of other reported systems such as the CaO–ZnO–SiO₂ [26] and CaO–ZnO–B₂O₃ [27] systems that were prepared in air.



Fig. 1. Phase diagram of the CaO–CoO_z–ZnO system at 885 °C in air, showing the limits of various solid solutions, and the tie-line relationships of various phases.

3.2. $CaO-CoO_z$ system

Two binary phases were determined in this system, namely, Ca₃Co₂O₆ and Ca₃Co₄O₉. The Ca₃Co₄O₉ phase has been extensively studied [28–32]. A very narrow solid solution region of $(Co,Ca)_3O_{4-x}$ was reported by Woermann and Muan [32] at around 890 °C. This region expands significantly at higher temperature. Ca₃Co₄O₉ is a misfit layered oxide that has two monoclinic subsystems with identical *a*, *c*, and β , but different *b* [29]. The first subsystem consists of triple rock-salt layers of Ca₂CoO₃ in the *ab*-plane while the second subsystem consists of a single CoO₂ layer, which has the CdI₂-type structure. This phase exhibits strong anisotropic thermoelectric properties in the *ab*-plane. It is thought that the increased scattering of phonons at the interface of misfit layers leads to the lowering of the lattice thermal conductivity. The chemical formula can be written as $[Ca_2CoO_3]RS[CoO_2]_{1.61}$, where RS is the rock salt and 1.61 expresses the incommensurable character for the *b* parameter of the rock salt and the CdI₂-type structure.

 $Ca_3Co_2O_6$ (*R*3*c*, *a*=9.0793(7), and *c*=10.381(1) Å [33]; Fig. 2) is the n=1 member of the homologous series with the general formula of $A_{n+2}B_nB'O_{3n+3}$. In this formula, A is an alkali-earth element such as Ca, Sr, and Ba [34], B describes the cobalt ion inside the octahedral cage, and B' is the cobalt ion inside a trigonal prism. $Ca_{n+2}Co_nCoO_{3n+3}$ consists of one-dimensional linear parallel Co₂O₆⁶⁻ chains, built by successive alternating face-sharing CoO₆ trigonal prisms and CoO₆ octahedra along the hexagonal *c*-axis [35]. This face-sharing feature is in contrast to $Ca_3Co_4O_9$ [12] and NaCo₂O₄ [8], which consist of edge-sharing CoO₆ octahedra. The linear $Co_2O_6^{6-}$ chains of $Ca_3Co_2O_6$ consist of one CoO₆ octahedron alternating with one CoO₆ trigonal prism. These $Co_2O_6^{6-}$ chains are separated by octa-coordinated Ca^{2+} ions. The compounds $A_{n+2}Co_nCo'O_{3n+3}$ can also be considered as ordered intergrowth between the $n = \infty$ (ACoO₃) and n = 1 $(A_3Co_2O_6)$ end members [36,37]. We found that when A = Ca, only the n=1 member, namely, Ca₃Co₂O₆, can be made.

3.3. $ZnO-CoO_x$ system

The only phases found in the $ZnO-CoO_x$ system are the two solid solutions formed in the vicinity of ZnO and CoO_x . Co substitutes into

the Zn site of ZnO ($P6_3mc$, a=3.250 Å, c=5.207 Å [38]) to form ($Zn_{1-x}Co_x$) O_{1-z} ($0 < x \le 0.06$)), and Zn substitutes into the tetrahedral Co site of Co_3O_4 ($Fd\bar{3}m$, a=8.066(3) Å [39]) to form ($Co_{1-x}Zn_x$)₃ O_{4-z} ($0 < x \le 0.17$, with the maximum substitution composition of ($Co_{2.49}Zn_{0.51}$) O_{4-z}). Apparently as temperature increases, the solid solution range of these two solid solutions increases as well. For example, it was reported that in the TiO₂–ZnO–CoO_z system [40], which was determined at 1050 °C in air, the solid solution limit



Fig. 2. Crystal structure of $Ca_3Co_2O_6$ (n=1 member in $Ca_{n+2}Co_nCoO_{3n+3}$), showing the feature of linear chains of successive alternations of CoO_6 octahedron with CoO_6 prism.

Table 2

Refined crystal structure of $(Zn_{0.98}Co_{0.02})O$; space group $P6_3mc$, a=3.2505(1)Å, c=5.2046 (1)Å, and V=47.623(1)Å³.

Atom	x	у	z	Site occ.	U _{iso (} Å ²)	Mult
Zn1	1/3	2/3	0.0	0.98	0.02600(19)	2
Co2	1/3	2/3	0.0	0.02	0.02600(19)	2
O3	1/3	2/3	0.3828(4)	1.0	0.0270(7)	2

of $(Zn_{1-x}Co_x)O_{1-z}$ was at a higher value of $x \approx 0.18$, and similarly, in $(Co_{3-x}Zn_x)O_{4-z}$, the limit was $x \approx 0.6$.

3.3.1. $(Zn_{1-x}Co_x)O(x=0.02)$

This green-colored phase crystallizes in the hexagonal wurtzite type structure in which an extensive network of cornershared ZnO₄ tetrahedra was found. Excellent Reitveld refinement results for $(Zn_{1-x}Co_x)O$ were obtained: $WR_p=0.0857$, $R_p=0.0675$, $\chi^2=3.391$, R(F)=0.0280, $R(F^2)=0.0586$. Table 2 gives the atomic coordinates and the distances. Essentially Co was found to dope in the Zn site. Fig. 3 gives the observed (crosses) and calculated (solid line) XRD intensities pattern at 295 K. The difference pattern is plotted at the same scale as the other patterns. The row of tick marks indicates the calculated peak positions. The refined Zn/Co site occupancy in the $(Zn_{1-x}Co_x)O$ phase is 0.98(2)/0.02. The refined space group and lattice parameters are $P6_3mc$, a=3.2505(1)Å, c=5.2046(1)Å, V=47.623(1)Å³, and Z=2. Fig. 4 gives the structure of $(Zn_{0.98}Co_{0.02})O$ projected along the *c*-axis. The metal-oxygen distances are 1.992(2)Å (\times 1) and 1.9733(6)Å



Fig. 4. Crystal structure of (Zn_{0.98}Co_{0.02})O (of hexagonal wurtzite type).



Fig. 3. Observed (crosses) and calculated (solid line) XRD intensities pattern for (Zn_{0.98}Co_{0.02})O at 295 K.

(\times 3); the six O–(Zn,Co)–O bond angles are 108.01(6) $^{\circ}$ (\times 3) and 110.95(5) $^{\circ}$ (\times 3).

The pattern of Fig. 3 was collected on a Bruker D2 Phaser diffractometer equipped with a LynxEye position-sensitive detector using CuK α radiation at a tube power of 30 kV and 10 mA. The pattern was collected from 4.99782–130.004° 2 θ in 0.0202144° steps, counting for 1 s/step.

3.3.2. Structure of $(Co_{1-x}Zn_x)_3O_{4-x}$ (x=0.3)

This black-colored phase has a cubic closed-packed spinelstructure. Fig. 5 gives the observed and calculated neutron powder diffraction intensities at room temperature. The refinement residues are: W R_p =4.92%, R_p =4.08%, and χ^2 =1.148.

In a normal AB_2O_4 spinel, A is a divalent atom and B is a trivalent atom. Table 3 gives the refined structure and atomic parameters for a selected representative composition of the normal spinel, $(Co_{2.7}Zn_{0.3})O_{4-z}$, at room temperature. $(Co_{2.7}Zn_{0.3})O_{4-z}$ is isostructural with Co_3O_4 , which crystallizes in the space group of $Fd\bar{3}m$ (a=8.08756(1) Å). In $(Co_{7}Zn_{3}O_{4-z}$, there are eight filled tetrahedral sites and 16 octahedral sites (8*a* and 16*d*) per unit cell. In one unit cell, there are 4 layers of BO₆ octahedral chains along the *c*-axis, with neighboring layers perpendicular to each other, as shown in Fig. 6. The CoO₄ tetrahedra are found between the CoO₆ octahedral chains (Fig. 7).

As the common coordination feature of Zn is tetrahedral, Zn was indeed found to substitute solely in the tetrahedral 8*a* site in the structure of $(Co_{2.7}Zn_{0.3})O_{4-z}$. The tetrahedral metal–oxygen distance in $(Co_{2.7}Zn_{0.3})O_{4-z}$ is 1.9434(3) Å (possibly of 2 + valence site), and the octahedral Co3–O5 distance is of an expected smaller magnitude of 1.9172(2) Å. It was reported that in a relatively highly doped spinel structure, $Co_{2.28}Cu_{0.72}O_4$ [41], Cu was found to substitute in both 8*a* and 16*d* sites of approximately equal quantities. Therefore it appears that when the dopant is of a



Fig. 5. Observed (crosses) and calculated (solid line) NPD intensities pattern for $(Co_{2,7}Zn_{0,3})O_{4-z}$ at 295 K.

Table 3

Refined structure and atomic parameters for Zn_{0.3}Co_{2.7}O_{4-z} at room temperature using neutron diffraction (space group: $Fd\bar{3}m$, a=8.08756(1) Å and V=528.99(2) Å³); $R_{wp}=4.92\%$ $R_p=4.08\%$, and $\chi^2=1.148$.

Atom	Frac	x	у	Z	100 U _{iso}
Co(8 <i>a</i>)	0.630(9)	1/8	1/8	1/8	$\begin{array}{c} 0.47(5) \\ 0.47(5) \\ 0.36(4) \\ 0.36(4) \\ 0.461(12) \end{array}$
Zn(8 <i>a</i>)	0.370(9)	1/8	1/8	1/8	
Co(16 <i>d</i>)	1.000(6)	1/2	1/2	1/2	
Zn(16 <i>d</i>)	0.000(6)	1/2	1/2	1/2	
O(32 <i>e</i>)	0.990(2)	0.26374(2)	0.26374(2)	0.26374(2)	



Fig. 6. Crystal structure of $(Co_{1-x}Zn_x)_3O_{4-z}$ showing neighboring octahedral layers perpendicular to each other.



Fig. 7. Crystal structure of $(Co_{1-x}Zn_x)_3O_{4-z}$ showing the arrangement of the CoO_4 tetrahedra.

small quantity, such as that in $(Co_{2.7}Zn_{0.3})O_{4-z}$, it will first enter the 8*a* site, and as the concentration of the dopant is increased, it would then enter the octahedral 16*d* site.

3.4. CaO-CoO_z-ZnO system

No new compound was found in the CaO–CoO_z–ZnO ternary system. Despite the report of a ternary Ca₃ZnCoO_{6-z} phase by Kawasaki and Takano [42], which is rhomobohedral, $R\bar{3}c$ (a=9.09 Å and c=10.61 Å) and crystallizes in the K₄CdCl₆ structure [43], this phase can only be prepared using the solution methods, followed by heat treatment under oxygen.

The tie-line relationships appear to be rather simple in the CaO-CoO_z-ZnO system. The limit of the two solid solution series, Ca₃(Co,Zn)₂O_{6-z} and Ca₃(Co,Zn)₄O_{9-z}, are Ca₃Co_{1.9}Zn_{0.1}O_{6-z} and Ca₃(Co,Zn)₄O_{9-z}. Tie-line bundles are found to be between Ca₃(Co,Zn)₄O_{9-z} and (Co,Zn)₃O₄, between Ca₃(Co,Zn)₂O_{6-z} and Ca₃(Co,Zn)₄O_{9-z}, and between Ca₃(Co,Zn)₂O_{6-z} and Ca₃(Co,Zn)₄O_{9-z}, between Ca₃(Co,Zn)₂O_{6-z} and Ca₃(Co,Zn)₄O₂ between Ca₃(Co,Zn)₂O_{6-z} and Ca₃(Co,Zn)₄O₂ between Ca₃(Co,Zn)₂O_{6-z} and Ca₃(Co,Zn)₄O₂ between Ca₃(Co,Zn)₄O₂ between Ca₃(Co,Zn)₄O₂ between Ca₃(Co,Zn)₂O_{6-z} and Ca₃(Co,Zn)₄O₂ between Ca₃(Co,Zn)O_{4-z} between Ca₃(Co,Zn)O_{4-z} between Ca₄(Co,Zn)₄O₄ between Ca₄(Co,Zn)₄O₄ between Ca₄(Co,Zn)₄O₄ between Ca₄(Co,Zn)₄O₄ between Ca₄(Co,Zn)₄O₄ between Ca₄(Co,Zn)O₄ between Ca₄(Co,Zn)O₄

The crystal structure of the Ca₃(Co,Zn)₂O_{6-z} phase that was determined by X-ray Rietveld refinement technique can be described by Fig. 2. The structure consists of linear chains of face-sharing (Co, Zn)O₆ trigonal prisms and CoO₆ octahedra. Table 4 gives the refinement results, atomic coordinates and bond distances for Ca₃(Co_{0.95}Zn_{0.05})₂O_{6-z}. As reported in literature, the Zn–O tetrahedral bond is about 1.98 Å [44], and the Zn–O and

Table 4

Refine	d cry	/stal stru	ictures of C	a ₃ (Co _{1.95} 2	Zn _{0.05})O _{6-z} ; s	pace
group	RЗc	(lattice	parameters	, atomic	coordinates,	and
bond o	listaı	nces).				

<i>a</i> (Å)	9.07881(1)
<i>c</i> (Å)	10.39449(1)
<i>V</i> (Å ³)	741.979(1)
Ca1, <i>x</i> 0¼	0.36898(2)
<i>U</i> ₁₁	0.00528(7)
U ₂₂	0.00555(10)
033	0.00371(10)
Co2/Zn3, 000	1/0
U_{11}	0.00281(8)
U ₃₃	0.00132(13)
Co4/Zn5, 00¼	
frac	0.95/0.05
U ₁₁	0.00638(8)
U ₃₃	0.00232(14)
06, <i>x</i>	0.17682(6)
y x	0.02463(6)
U _{iso}	0.00530(12)
$C_{21} = C_{21} = C$	2 3465(5)
~ 2 (Å)	2.4657(5)
$\times 2$ (Å)	2.5596(5)
$\times 2$ (Å)	2.4617(5)
	1 0111(5)
$C02-U6 \times 6$ (A) O6-C02-O6 (deg.)	86.07(2)
00 c02 00 (acg.)	93.93(2)
$Co4 O6 \times 6$ (prismatic) (Å)	2.0713(5)
06-Co4-O6 (deg.)	78.05(2)
	87.72(3)
	127.60(3)
	147.53(3)
R _{wp}	0.0975
$\frac{R_{\rm p}}{\gamma^2}$	0.0777
$\frac{\lambda}{R(F)}$	0.0339
$R(F^2)$	0.0426
$\Delta F +$	5.3
$\Delta F -$	-2.1
Profile coefficients	2 22(4)
U W	2.33(4) 0.0655(5)
X	0.233(3)
ptec	-0.173
S(400)	0.00096(3)
S(004) S(202)	0.00358(4) 0.00133(3)
3(232)	



Fig. 8. EXFAS data for Ca₃(Co_{0.95}Zn_{0.05})₂O_{6-z}.

Co-O octahedral bond lengths are at longer values of 2.083(3) and 2.0767(8) Å, respectively [45]. The prismatic Zn–O and Co–O bond distances are at further longer distances of 2.199(4) Å [46] and 2.018 Å [47], respectively. In the present structure, the average octahedral Co-O distance is 1.9111(5)Å, whereas the average Co-O prismatic distance is longer, 2.0713 (5) Å; therefore Zn most likely substitutes onto the prismatic Co-O site. However, as mentioned before, the X-ray diffraction technique is a technique for obtaining average Zn-O and Co-O bond distances and cannot conclusively provide the local structure information. The EXAFS results that provide local bonding environments are shown in Fig. 8. The white lines for the zinc nitrate and $Ca_3Co_{1.95}Zn_{0.05}O_{6-z}$ are almost at the same position, which is a signature of zinc being in 6-coordination state. The small pre-edge peak coincides with that of Zn in ZnO. This would suggest that Zn exists in a small amount in tetrahedral form (possibly in the amorphous form on the surface as a second phase). Zn has a small fraction residing in the tetrahedral site with a bond length 1.89 ± 0.05 Å and the main concentration is in the 6-coordination site with 2.11(2) Å (which is longer than the average octahedral Zn-O distance of 1.9111(5) Å). The 2.11(2) Å of Zn–O would suggest a bond length of a trigonal prism environment.

3.5. Thermoelectric properties

The thermoelectric properties of a number of compounds in the Ca–Co–Zn–O system have been reported in literature [33,48–55]. In the Ca–Co–O system, the Ca₃Co₄O₉ phase was reported to exhibit strong anisotropic properties and has excellent thermoelectric property in the *ab*-plane. *ZT* was reported to be ≈ 1 at 1000 K. It is thought that the increased scattering of phonons at the interface of misfit layers leads to the lowering of the lattice thermal conductivity. The Seebeck coefficient for single crystal Ca₃Co₂O₆ has been reported by Mikami and Funahashi [10] to be relatively high and positive, and the thermal conductivity is relatively low at high temperature. The transport properties are dominated mainly by p-type carriers. *ZT* was determined to be about 0.15 at 1000 K for a single crystal of Ca₃Co₂O₆.

The Seebeck coefficient, thermal conductivity, resistivity, and *ZT* values of Ca₃Co₂O₆, Ca₃(Co, Zn)₂O_{6-z}, Ca₃(Co, Zn)₄O_{9-z}, and Ca₃(Co, Zn)₄O_{9-z} as a function of temperature, from this study, are given in Fig. 9a–d. The transport properties are dominated mainly by p-type carriers, with high temperature magnitudes between 200 μ V/K at 900 K and 400–700 μ V/K at 300 K for the



Fig. 9. (a) Seebeck coefficient, (b) resistivity, (c) thermal conductivity and (d) figure of merit, ZT, of $Ca_3Co_2O_6$, $Ca_3(Co, Zn)_2O_{6-z}$, $Ca_3Co_4O_9$, and $Ca_3(Co, Zn)_4O_{9-z}$ as a function of temperature.

Ca₃(Co,Zn)₂O_{6-*z*} phase, and about 190 μ V/K between 300 and 900 K for the Ca₃(Co,Zn)₄O_{9-*z*} phase. The thermal conductivity data of the two solid solution series are similar and are reasonably low in the temperature range of 300–900 K. Zn substitution does not seem to have much effect on the thermoelectric properties of the Ca₃(Co,Zn)₄O_{9-*z*} series. Although Zn-doping lowers the resistivity of the Ca₃(Co,Zn)₂O_{6-*z*} series between 300 and 600 K, probably due to the changes of carrier concentration, at 900 K the values converge. The *ZT* value for our Ca₃Co₄O₉ sample (Fig. 9d) is not as high as reported [48], which is probably due to the fact that the processing conditions of our sample have not been optimized. The Ca₃(Co,Zn)₄O_{9-*z*} series has a relatively higher *ZT* values than that of the Ca₃(Co,Zn)₂O_{6-*z*} series. At present, the Ca₃Co₄O₉ compound remains the material that gives the highest *ZT* values among the oxides studied [48].

4. Summary

The phase diagram of the Ca–Co–Zn–O system that was determined at 885 °C in air provides detailed compatibility relationships in the binary and the ternary oxide systems that are important for processing and for the understanding of material properties. The Ca–Co–Zn–O phase diagram consists of two low-dimensional phases, namely, Ca₃(Co,Zn)₄O_{9–z} with misfit layered structure, and Ca_{n+2}(Co,Zn)_n(Co,Zn)'O_{3n+3–z} with one-dimensional cobalt oxide chains that offer interesting thermoelectric properties. The homogeneity ranges of the Ca₃(Co,Zn)₄O_{9–x} and Ca_{n+2}(Co,Zn)_n(Co,Zn)'O_{3n+3–z} solid solutions have been determined. Zn substitution does not have much effect on the thermoelectric properties of the Ca₃(Co,Zn)₄O_{9–z}

series. Although the Ca_{n+2}(Co,Zn)_n(Co,Zn)'O_{3n+3-z} oxides exhibit relatively high Seebeck coefficient and low thermal conductivity in general, their resistivity values, however, are relatively high. Although Zn-doping lowers the resistivity of the Ca₃(Co,Zn)₂O_{6-z} series between 300 and 600 K, the Seebeck coefficients, however, also become lower. Unless we can decrease the resistivity of the Ca₃Co₂O₆ type phases via proper substitution and improved processing, the best cobaltite materials for thermoelectric applications, at present, remain to have the Ca₃Co₄O₉ structure.

Acknowledgments

Use of the Advanced Photon Source at Argonne National Laboratory was supported by the U. S. Department of Energy, Office of Science, Office of Basic Energy Sciences, under Contract no. DE-AC02-06CH11357. One of the authors (SC) would like to thank Dr. J.T. Miller (ANL) and Dr. T. Shibata (IIT) for useful discussions.

References

- G.S. Nolas, J. sharp, H.J. Goldsmid, Thermoelectric: Basic Principles and New Materials Developments, Springer, New York, 2001.
- [2] T.M. Tritt, M.A. Subramanian, guest editors, Harvesting energy through thermoelectrics: power generation and cooling, MRS Bulletin, Materials Research Society (2006).
- [3] T.M. Tritt, Science 272 (1996) 1276.
- [4] Kuei Fang Hsu, Sim Loo, Fu Guo, Wei Chen, Jeffrey S. Dyck, Ctirad Uher, Tim Hogan, E.K. Polychroniadis, Mercouri G. Kanatzidis, Science 303 (2004) 818.
- [5] R. Venkatasubramanian, E. Siivola, T. Colpitts, B. O'Quinn, Nature 413 (2001) 597.

- [6] S. Ghamaty, N.B. Eisner, Proceeding of Interpack 2005: ASME Technical Conference on Packaging of MEMS, NEWS and Electric Systems, July 17–22, 2005, San Francisco, CA.
- [7] M.S. Dresselhaus, G. Chen, M.Y. Tang, R.G. Yang, H. Lee, D.Z. Wang, Z.F. Ren, J.P. Fleurial, P. Gogna, Materials Research Society Symposium Proceedings 886, Materials and Technologies for Direct Thermal-to-Electric Energy Conversion, 2006, pp. 3–12.
- [8] I. Terasaki, Y. Sasago, K. Uchinokura, Phys. Rev. B 56 (1997) 12685-12687.
- [9] M. Mikami, R. Funashashi, M. Yoshimura, Y. Mori, T. Sasaki, J. Appl. Phys. 94 (10) (2003) 6579–6582.
- [10] M. Mikami, R. Funahashi, J. Solid State Chem. 178 (2005) 1670-1674.
- [11] D. Grebille, S. Lambert, F. Bouree, V. Petricek, J. Appl. Crystallogr. 37 (2004) 823-831.
- [12] A.C. Masset, C. Michel, A. Maignan, M. Hervieu, O. Toulemonde, F. Studer, B. Raveau, Phys. Rev. B 62 (2000) 166–175.
- [13] H. Minami, K. Itaka, H. Kawaji, Q.J. Wang, H. Koinuma, M. Lippmaa, Appl. Surf. Sci. 197 (2002) 442–447.
- [14] Y.F. Hu, W.D. Si, E. Sutter, Q. Li, Appl. Phys. Lett. 86 (2005) 082103.
- [15] Powder Diffraction File (PDF), produced by ICDD, Newtown Square, 12
- Campus Blvd., Newtown Squares, PA 19073-3273, USA. [16] J. Wang, B.H. Toby, P.L. Lee, L. Ribaud, S. Antao, C. Kurtz, M. Ramanathan,
- R.B. Von Dreele, M.A. Beno, Rev. Sci. Instrum. 79 (2008) 085105. [17] P.L. Lee, D. Shu, M. Ramanathan, C. Preissner, J. Wang, M.A. Beno, R.B. Von
- Dreele, L. Ribaud, C. Kurtz, S.M. Antao, X. Jiao, B.H. Toby, J. Synchr. Radiat. 15 (2008) 427–432.
 [18] C. Preissner, D. Shu, B.H. Toby, P. Lee, J. Wang, D. Kline, K. Goetze, Nucl.
- Instrum. Meth. Phys. Res. A, in press.
- [19] L.R. Dalesio, J.O. Hill, M. Kraimer, S. Lewis, D. Murray, S. Hunt, S.W. Watson, M. Clausen, J. Dalesio, Nucl. Instrum. Meth. Phys. Res. A 352 (1994) 179–184.
- [20] Standard reference materials (SRMTM) are produced by National Institute of Standards SRM Office, Gaithersburg, MD 20899. For details, please contact srminfo@nist.gov.
- [21] H.M. Rietveld, J Appl. Cryst. 2 (1969) 65-71.
- [22] A.C. Larson, R.B. von Dreele, GSAS—General Structure Analysis System, US Government contract (W-7405-ENG-36) by the Los Alamos National laboratory, which is operated by the University of California for the U.S. Department of Energy, 1992.
- [23] C.U. Segre, et al., Synchrotron Radiation Instrumentation: Eleventh U.S. Conference CP521, 2000, pp. 419–422.
- [24] M. Newville, J. Synchr. Radiat. 8 (2001) 322-326.
- [25] B. Ravel, M. Newville, J. Synchr. Radiat 12 (4) (2005) 537-541.
- [26] E.R. Segnit, J. Am. Ceram. Soc. 37 (6) (1954) 273-377.

- [27] A. Whitaker, D.F. Borton, J. Mater. Sci. 7 (1972) 194-197.
- [28] D. Grebille, S. Lambert, F. Bouree, V. Petricek, J. Appl. Crystallogr 37 (2004) 823–831.
- [29] A.C. Masset, C. Michel, A. Maignan, M. Hervieu, O. Toulemonde, F. Studer, B. Raveau, Phys. Rev. B 62 (2000) 166–175.
- [30] H. Minami, K. Itaka, H. Kawaji, Q.J. Wang, H. Koinuma, M. Lippmaa, Appl. Surf. Sci. 197 (2002) 442–447.
- [31] Y.F. Hu, W.D. Si, E. Sutter, Q. Li, Appl. Phys. Lett. 86 (2005) 082103.
- [32] E. Woermann, A. Muan, J. Inorg. Chem 32 (1970) 1455.
- [33] H. Fjellvag, E. Gulbrandsen, S. Aasland, A. Olsem, B.C. Hauback, J. Solid State Chem. 124 (1996) 190.
- [34] W. Wong-Ng, G. Liu, J. Martin, E. Thomas, N. Lowhorn, M. Otani, J. Appl. Phys. 107 (2010) 033508.
- [35] T. Takami, H. Ikuta, U. Mizutani, Jpn. J. Appl. Phys. 43 (22) (2004) 8208–8212.
 [36] K. Boulahya, M. Parras, J.M. González-Calbet, J. Solid State Chem. 142 (1999)
- 419-427.
- [37] K. Boulahya, M. Parras, J.M. González-Calbet, J. Solid State Chem. 145 (1999) 116-127.
- [38] H. McMurdie, M. Morris, E. Evans, B. Paretzkin, W. Wong-Ng, L. Ettinger, C.R. Hubard, Powder Diff. 1 (1988) 76.
- [39] W.L. Roth, J. Phys. Chem. Solids 25 (1964) 1.
- [40] A. Navrotsky, A. Muan, J. Inorg. Nucl. Chem. 32 (11) (1970) 3471-3484.
- [41] K. Krezhov, K. Petrov, T. Karamaneva, J. Solid State Chem. 48 (1983) 33-39.
- [42] S. Kawasaki, M. Takano, T. Inami, J. Solid State Chem. 145 (1999) 302-308.
- [43] G. Bergerhoff, D.O. Schmitz, Z. Anorg. Allg. Chem. 284 (1956) 10.
- [44] C. Tusche, H.L. Meyerheim, J. Kirschner, Phy. Rev. Lett 99 (2007) 026102.
- [45] J.Y. Lu, M.A. Lawandy, J. Li, T. Yuen, C.L. Lin, Inorg. Chem. 38 (1999) 2695.
- [46] Lampe-önnerud, H.-C. zur Loye, Inorg. Chem. 35 (1996) 2155-2156.
- [47] W. Wong-Ng, J.A. Kaduk, G. Liu, Powder Diff. 26 (2011) 22.
- [48] M. Mikami, R. Funahashi, IEEE Proceedings on 22nd International Conference on Thermoelectrics, 2003, pp. 200–202.
- [49] S. Hébert, D. Flahaut, C. Martin, S. Lemonnier, J. Noudem, C. Goupil, A. Maifnan, J. Hejtmanek, Prog. Solid State Chem 35 (2007) 457-467.
- [50] K. Iwasaki, H. Yamane, S. Kubota, J. Takahashi, M. Shimada, J. Alloy. Compd. 358 (2003) 210–215.
- [51] Y. Miyazaki, Solid State Ionics 172 (2004) 463-467.
- [52] S. Li, R. Funahashi, I. Matsubara, H. Yamada, K. Ueno, S. Sodeoka, Ceram. Int. 27 (2001) 321–324.
- [53] S. Li, R. Funahashi, I. Matsubara, H. Yamada, K. Ueno, S. Sodeoka, H. Yamada, Chem. Mater. 12 (2000) 2424–2427.
- [54] M. Prevel, O. Perez, J.G. Noudem, Solid State Sci. 9 (2007) 231-235.
- [55] T. Takama, H. Ikuta, U. Mizutani, Jpn. J. Appl. Phys. 43 (12) (2004) 8208–8212.